

AN EPITOME OF CONFIGURATIONAL DATA ON CONNECTED CLUSTERS

J. A. M. S. DUARTE

Departamento de Física — Faculdade de Ciências do Porto, 4000 Porto, Portugal

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ABSTRACT — New lattice data on configurational histograms are given for bond and site clusters grouped by fixed percolation perimeter, fixed energy perimeter and fixed cluster size. The latter are illustrated by several combinations of interest of cyclomatic number discriminations.

INTRODUCTION

It has long been recognized that configurational studies are a fundamental tool in the theory of critical phenomena. Recently, however, powerful techniques (like transfer-matrix renormalization and field theoretical methods [17]) have surged on to the statistics of lattice clusters in the percolation and animal problems (see e. g. ref [18]) and significant advances in the knowledge of the critical exponents for both problems have been brought close to a virtually «exact» solution. There is, however, still open a rich field of specializations (valence, cyclomatic number, specific connectivity requirements, restricted sets of clusters (animals) defined through topological constraints). Our aim in this paper is twofold: written in mid-81 it should concentrate on selected topics referring to the cluster topology which are likely to assume physical relevance in the future, and where series expansions and configurational studies will remain competitive. On the other hand, it should unify various treatments that have remained scattered in the literature without any systematic exploration (like bond or site content in percolation). We have divided the data in 5 broad groups: fixed energy groupings, fixed percolation perim-

eter groupings, fixed size percolation groupings, cyclomatic number distributions (in percolation) and fixed size energy groupings. Each one of them is preceded by a succinct description of the graph theoretical procedures used in its derivation.

The notation to be consistently applied throughout the paper is:

- s — denotes the number of cluster sites
- b — denotes the number of cluster bonds
- c = b - s + 1 denotes the cyclomatic number of a connected cluster
- e — denotes the external bond (energy) perimeter
- t — denotes the perimeter in the percolation sense

g_{se} ; g_{sbt} — give the number of geometrically different cluster configurations with a given label s,e or s,b,t .

Note that the normalization of the various $g_{...}$ may occasionally vary for convenience. We have indicated in each case the factor relative to a normalization per lattice site.

A — Fixed energy groupings

Whenever the bond perimeter of connected site clusters is fixed, the resulting distributions according to the variable number of sites enclosed within a given configuration of boundary bonds can easily be translated topologically into a fixed perimeter — enclosed area problem by considering the dual lattice (Sykes et al [1], [2]). Consider figure 1 for the triangular — honeycomb system: in Fig. 1 A, the connected cluster of 8 sites and 11 bonds on the triangular lattice is the *dual* of the honeycomb configuration with 26 sides and area 8. Denoting the number of sites by s , the number of (*internal*) bonds by b and external bonds by e , the following linkage rule for site clusters (strongly embedded clusters)

$$e = zs - 2b$$

A. 1

is valid on any lattice (coordination number z). For configurations of the type in Fig. 1A there are no sites enclosed within

the configuration and not belonging to it but Fig. 1B indicates the possibility of such configurations: all clusters that can be derived from the fully compact cluster limited by the outermost boundary through the exclusion of any combination of hexagonal faces marked with (x) are still duals of connected clusters on the triangular lattice, but, unlike the case of Fig. 1A, their boundary is no longer singly-connected (it is no longer a simple polygon).

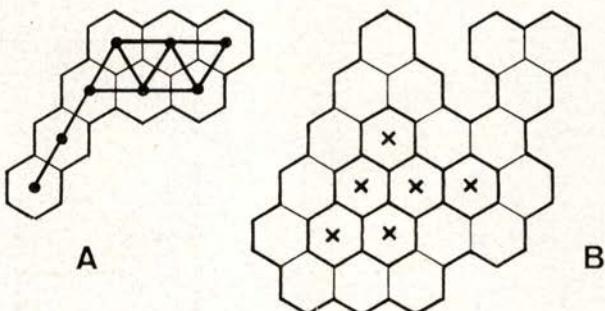


Fig 1

- A — Site cluster on the triangular lattice and its dual on the honeycomb lattice. The triangular configuration is compact (no inner perimeter sites) and its dual is bounded by a simple polygon.
- B — Another example of a honeycomb configuration. Exclusion of any face marked with (x) generates a connected dual from the larger simple polygon.

The situation recurs for the simple quadratic lattice, which is well known to be self-dual (Fig. 2). Fig. 2C is a compact configuration bounded by a simple polygon (it is, in fact, the isoperimetric solution for perimeter 18, Duarte and Marques [3] — and area 20). Once again, exclusion of any combination of square faces marked with (x) generates a connected area (alternative examples are drawn in 2A and 2B), which is still a dual of some site cluster on the same lattice. Fig. 2D shows, explicitly, a square site tree (17 sites) and its connected dual.

Now all perimeter distributions of site clusters contribute to the low temperature ferromagnetic polynomials for the Ising

model [1]. It is, however, required, for their isolation, to separate the contribution of multicomponent graphs as described in [4]. To go one stage further and separate the simple polygons from the nonpolygonal connected duals, we note that on the honeycomb lattice (Fig. 1) it is impossible to have more than 3 ele-

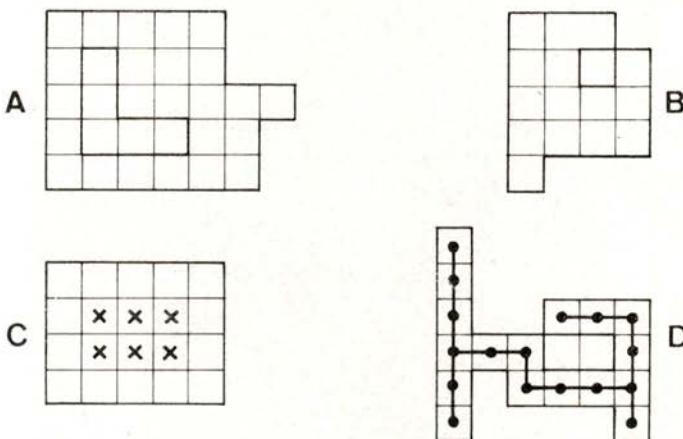


Fig. 2

A, B — Examples of connected duals on the square lattice.

C — Simple polygon on the square lattice. Exclusion of any face marked with (x) generates a connected dual from the larger configuration.

D — A square lattice tree and its dual.

mentary hexagonal faces meeting at a site and, therefore, the contribution from those configurations can be singled out from clusters discriminations on the triangular lattice taking into account the number of elementary triangular faces f (as well as s and b). Isolated inner boundary sites will then occur for all clusters where f does not account for the total cyclomatic number:

$$b - s + 1 \neq f \quad \text{A. 2}$$

and this inequality identifies the non-polygonal connected duals: all inner boundaries will be separate from the outermost boundary.

It is impossible to proceed in this way for the square lattice: a tree like in Fig. 2D does not verify A.2 and yet its dual is not a polygon. In general the problem of determining the connected duals up to a reasonable order is lessened by simple conversion of fixed s distributions (b groupings), like those below in section D. Earlier results for the square polygons can be found in Hiley and Sykes [5].

The additional data should sum to the known results for the total number of polygons (fixed perimeter), greatly extended in a recent paper by Enting [6] through the use of transfer matrix techniques.

We give new data for polygons on the honeycomb ($e \leq 42$) and square lattices ($e \leq 22$) as well as for the corresponding connected duals.

B — Fixed percolation perimeter groupings

When the perimeter is measured in the percolation sense, i.e. by the number of sites (bonds) t necessary for the isolation of a given cluster on a lattice, the perimeter groupings suffer considerable rearrangement of the various cluster contributions. The usual perimeter method can, of course, be used for obtaining these groupings — once again, they can be obtained through a straightforward conversion of the fixed size percolation distributions, although such information must be completed (for detailed descriptions see Sykes et al [4], Blease et al [7]).

In this paper we present results for these groupings on the square ($t \leq 16$) and honeycomb ($t \leq 13$) lattices (site problem) as well as for the Kagomé site problem ($t \leq 16$). As in section A, the problem is equivalent to the enumeration of the histograms g_{st} or g_{bt} at fixed t ; g_{st} or g_{bt} gives the number of geometrically different cluster configurations per site (or per bond) with a given perimeter value t (here t refers to bond and site perimeter for bond and site percolation respectively).

In addition, we have used inequality A.2 (and further discrimination through b , s , t and f) to isolate all the non-polygonal connected duals of the triangular lattice according to their percolation weight. The resulting perimeter polynomials are given

through order 21 (they should be compared with the complete set of perimeter groupings for the problem, given in Sykes et al [4] ($t \leq 22$)).

C — Fixed size percolation groupings

The g_{st} (fixed size s) are the best illustrated groupings in the literature. They have been listed for 2, 3 and higher dimensions [8], [9], for both site and bond [10] problems on most usual lattices. The interested reader should refer to those papers for an outline of the method and detailed considerations on the applicability of the corresponding series expansions. We have added in this paper the groupings for the site problem on the 2 archimedean lattices of coordination number 5, (3,3,3,4,4) and (3,3,4,3,4) ($s \leq 12$). Both lattices (their Ising points are known exactly) provide good testing ground for the variation of the perimeter — to — size ratio and its connection with criticality (Duarte [11]). The well known sum rule to be verified by the g_{st} is

$$p = \sum_{s,t} s g_{st} p^s (1-p)^t \quad C. 1$$

D — Cyclomatic number distributions in percolation

A different type of configurational weighting which has been the object of much recent interest is the set of three — indexed discriminations of clusters by their site, bond content and perimeter (in the percolation sense). Through Euler's law these discriminations lead to the expansions of the average cyclomatic number $\langle c \rangle$ (Cherry [12], Gaunt et al [13]).

$$\langle c \rangle = \langle b \rangle - \langle s \rangle + \langle 1 \rangle \quad D. 1$$

and from expansion of the higher moments of the cluster size distribution of the type

$$\langle s^k \rangle = \sum_{s,b,t} s^k g_{sbt} p^b (1-p)^t \quad D. 2$$

for bond percolation and

$$\langle b^k \rangle = \sum_{s,b,t} b^k g_{sbt} p^s (1-p)^t \quad D. 3$$

for site percolation, new quantities of interest, paralleling the moments in the usual cluster size distribution are obtained. They are expected to belong to the same universality class as normal percolation and therefore constitute alternative ways of calculating the critical exponents for percolation. For $k=2$, D.2 and D.3 lead to «susceptibility» series diverging near p_c with a critical exponent γ , like the mean cluster size

$$p \rightarrow p_c^-, \langle s^2 \rangle_{\text{bond}} \sim |p_c - p|^{-\gamma}, \langle b^2 \rangle_{\text{site}} \sim |p_c - p|^{-\gamma}, \quad D. 4$$

This property has been occasionally used in the literature (Dunn et al [14], Agrawal et al [15]). A systematic study for 2 dimensional percolation (p_c lattices) is reported in [16].

We present results for the set of histograms $\Sigma_b b g_{sbt}$ for the triangular, square matching, Kagomé, honeycomb and archimedean (3,3,3,4,4) and (3,3,4,3,4) site problems and for $\Sigma_s s g_{sbt}$ for the square and honeycomb bond problems, as well as for $\Sigma_b b^2 g_{sbt}$ for the Kagomé site problem. We recall that for the first moment distributions the sum rules

$$\sum_{b,s,t} b g_{sbt} p^s (1-p)^t = (z/2) p^2 \quad D. 5$$

for site percolation and

$$\sum_{b,s,t} s g_{sbt} p^b (1-p)^t = 1 - (1-p)^z \quad D. 6$$

for bond percolation, should be verified.

It also seems adequate to mention that the use of detailed valence discriminations constitute an alternative way of determining their cyclomatic number distributions. Since they represent an expansion from 3-indexed to z-indexed discriminations it is usually more cumbersome to take this line of procedure (it was however followed in refs [12], [13]). If the sites in a connected cluster are partitioned according to the number of

site neighbours in the cluster (their valence) the following linkage rules are verified

$$\sum_v s_v = s \quad D. 7$$

$$\sum_v v s_v = 2 b \quad D. 8$$

with s_v the number of sites with valence v and where the summations run from 1 to z . Equations D. 7 and D. 8 are unwieldy for bond percolation (valence considerations apply equally to bond and site percolation clusters), since b is usually fixed and any cyclomatic number mixing in a given g_{bt} must be disentangled from a combination of different s as well as from all compatible combinations of s_v 's. It is generally more direct to start from bond percolation distributions and exploit the following properties of the yield factor generation (see Bleasdale et al [7] for an exposition of the method):

- a) For a given space-type strongly embeddable on the specific lattice under investigation, the number of bonds that can be transferred from the bond content to the bond perimeter is not greater than the cyclomatic number, if connectivity is to be preserved.
- b) The number of transferable bonds is zero for strongly embedded trees.
- c) The length of the bond tree percolation polynomial (with $b = s - 1$ bonds) is $c_{\max} + 1$, where c_{\max} is the maximum cyclomatic number of s -site clusters.
- d) Strongly embedded clusters always maximise the bond perimeter for given values of b and c . Recalling that the linkage rule for strongly embedded clusters is $e = t = zs - 2b$, it can be seen that the difference in bond perimeter between clusters with successive cyclomatic numbers (same b) is z . These properties enable a separation of the cyclomatic number contributions in bond perimeter polynomials of not too high b (like those in Sykes et al [10]). In every case the sum rule D. 6 acts as a check on such graph theoretical manipulations and use of valence discriminations is thereby avoided.

E — Fixed size energy groupings

These enumerations are converse of those considered in section A. For site clusters, these groupings constitute (through eq. A. 1) a strict partition according to the cyclomatic number (or alternatively, according to the number of bonds b in the cluster). No data exist in the literature regarding this specific partition, which has emerged in recent times as a very relevant tool for the study of branched polymers in the dilute limit (mainly through the studies of Lubensky and coworkers [17], [18]). Clearly, in this partition the 3-indexed discriminations in D. 2 and D. 3, the g_{sbt} , are summed over the perimeter index, so that with the resulting g_{sb} (which are equivalent to the g_{se}) new moments of the «animal distribution» [17] can be defined and numerically investigated.

We present data on the three 2-dimensional regular lattices. The first noticeable difference with respect to the g_{ts} is that the corresponding histograms evolve very slowly in shape, so that it might be argued that for this specific partition the lattices are not very effectively sampled. Now, in each case, the maximum bond perimeter corresponds to the minimum number of bonds in the cluster, so that the last value in each of the g_{se} histograms just gives the total number of site trees on each lattice. As we have mentioned, in the previous section, this total number will also appear as the maximum perimeter configuration value in the bond percolation polynomial of bond size $b = s - 1$. Hence, the present data represent an extension over the data of ref [10].

Unfortunately the following term in g_{se} (corresponding to the total number of polygons, tadpoles and other configurations of cyclomatic number 1) does not grow sufficiently fast for a non-degenerate histogram to occur for loose-packed lattices. Consideration of the bond case only worsens the balance of the histogram: A. 1 is no longer valid, so that the g_{sb} with $s = b + 1$, gives the total number of bond trees and through the use of the yield factor generation all site clusters of $b + 1$ sites give non-zero contributions to $g_{b+1,b}$.

In order to avoid these problems one must concentrate on high coordination number (site) lattices where the strong embeddability «propagates» the distributions towards lower e values. But,

better still, the careful exploitation of local linkage rules and consideration of both site and bond valence and the changes they undergo in bond-to-site transformations have enabled Duarte and Ruskin [19] to identify the valence structure on lattices which are covering lattices of bond problems. For the site trees ($g_{b+1,b}$) are always neighbour avoiding walks, belonging to a totally different universality class from branched trees and with a comparatively smaller growth parameter. The same happens for the terms of the form g_{bb} which originate from the corresponding bond polygons and bond trees with one single site of valence 3. Hence the g_{se} for the corresponding site covering problem shows a rapid evolution towards non-degenerate configurational histograms. We illustrate this point with the square covering $g_{se}(s \leq 13)$.

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APPENDIX

FIXED ENERGY GROUPINGS

A — Honeycomb polygons

	$e = 6$	$g_{es} (\times 1)$	$e = 20$	
1		1	5	60
	$e = 10$		6	42
2		3	7	30
	$e = 12$		8	6
3		2		$e = 22$
	$e = 14$		5	99
3		9	6	129
4		3	7	105
	$e = 16$		8	69
4		12	9	27
5		6	10	3
	$e = 18$			$e = 24$
4		29	6	280
5		21	7	276
6		14	8	246
7		1	9	160

10	86	17	1368
11	24	18	606
12	2	19	198
	e = 26	20	42
6	348	21	6
7	726		e = 34
8	720	8	4644
9	609	9	18786
10	432	10	27627
11	249	11	33405
12	117	12	32061
13	27	13	29097
14	3	14	23553
	e = 28	15	18597
7	1242	16	13128
8	1710	17	8877
9	1812	18	5412
10	1458	19	2943
11	1164	20	1401
12	702	21	507
13	414	22	147
14	168	23	27
15	42	24	3
16	6		e = 36
	e = 30	9	23472
7	1260	10	54148
8	3759	11	76662
9	4611	12	88378
10	4769	13	86860
11	3870	14	78978
12	3163	15	67134
13	2126	16	53826
14	1320	17	40866
15	729	18	29076
16	290	19	19672
17	87	20	12006
18	14	21	6936
19	1	22	3424
	e = 32	23	1458
8	5436	24	496
9	9804	25	128
10	12186	26	24
11	12030	27	2
12	10476		e = 38
13	8406	9	17382
14	6336	10	90924
15	4134	11	160131
16	2622	12	218436

13	240579	27	13128
14	242613	28	6060
15	219816	29	2412
16	193602	30	798
17	158682	31	216
18	126099	32	42
19	93948	33	6
20	68019	e = 42	
21	45531	10	65822
22	29049	11	431448
23	17115	12	897289
24	9138	13	1369834
25	4338	14	1691994
26	1719	15	1865164
27	579	16	1873893
28	147	17	1778925
29	27	18	1601354
30	3	19	1397388
e = 40		20	1168533
10	100740	21	951897
11	287838	22	742157
12	464580	23	564297
13	604434	24	410122
14	661206	25	288397
15	669792	26	192099
16	619944	27	122932
17	553584	28	73674
18	469290	29	41040
19	384144	30	21083
20	300192	31	9632
21	226296	32	3918
22	163500	33	1341
23	111960	34	392
24	73266	35	87
25	44646	36	14
26	25626	37	1

A — Square polygons

e = 4		9	656
e = 6		10	482
e = 8		11	310
e = 10	(see Hiley, Sykes [5])	12	151
e = 12		13	68
e = 14		14	22
e = 16	g _{se} ($\times 1$)	15	6
7	566	16	1
8	676		

e = 18		21	187
8	1868	22	68
9	2672	23	22
10	2992	24	6
11	2592	25	1
12	2086		e = 22
13	1392	10	21050
14	864	11	39824
15	456	12	56162
16	218	13	61032
17	88	14	60864
18	30	15	54032
19	8	16	45936
20	2	17	35952
		18	26858
e = 20		19	18744
9	6237	20	12456
10	10376	21	7648
11	13160	22	4472
12	12862	23	2408
13	11717	24	1208
14	9332	25	560
15	7032	26	238
16	4748	27	88
17	3010	28	30
18	1728	29	8
19	914	30	2

A — Honeycomb duals

e = 6		e = 26	
e = 10		6	348
e = 12		7	732
e = 14		8	720
e = 16		9	609
e = 18		10	432
e = 20		11	249
e = 22		12	117
e = 24	g _{se} ($\times 1$)	13	27
		14	3
6	281		e = 28
7	276	7	1248
8	246	8	1737
9	160	9	1818
10	86	10	1458
11	24	11	1164
12	2	12	702
		13	414

14	168	23	27
15	42	24	3
16	6	e = 36	
	e = 30	9	23662
7	1260	10	54411
8	3795	11	78990
9	4697	12	91389
10	4817	13	89680
11	3876	14	81183
12	3163	15	68294
13	2126	16	54261
14	1320	17	40950
15	729	18	29083
16	290	19	19672
17	87	20	12006
18	14	21	6936
19	2	22	3424
	e = 32	23	1458
8	5472	24	496
9	9990	25	128
10	12453	26	24
11	12264	27	2
12	10557	e = 38	
13	8418	9	17382
14	6336	10	92205
15	4134	11	165144
16	2622	12	226125
17	1368	13	250641
18	606	14	252177
19	198	15	227970
20	42	16	199110
21	6	17	161712
	e = 34	18	127287
8	4644	19	94242
9	19014	20	68067
10	28305	21	45531
11	34263	22	29049
12	32901	23	17115
13	29601	24	9138
14	23721	25	4338
15	18627	26	1719
16	13128	27	579
17	8877	28	147
18	5412	29	27
19	2943	30	3
20	1401	e = 40	
21	507	10	101679
22	147	11	295356

12	482850	13	1431751
13	631218	14	1782600
14	694122	15	1971774
15	702816	16	1985091
16	648951	17	1879842
17	575856	18	1685370
18	484095	19	1458954
19	392556	20	1210437
20	303741	21	975901
21	227472	22	753821
22	163743	23	568761
23	111990	24	411402
24	73266	25	288661
25	44646	26	192123
26	25626	27	122932
27	13128	28	73674
28	6060	29	41040
29	2412	30	21083
30	798	31	9632
31	216	32	3918
32	42	33	1341
33	6	34	392
e = 42		35	87
10	65822	36	14
11	438264	37	1
12	929414		

A — Square duals

e = 4	g _{se} ($\times 1$)	8	134
1	1	9	72
e = 6		10	30
2	2	11	8
e = 8		12	2
3	6	e = 16	
4	1	7	570
e = 10		8	677
4	18	9	656
5	8	10	482
6	2	11	310
e = 12		12	151
5	55	13	68
6	40	14	22
7	22	15	6
8	6	16	1
9	1	e = 18	
e = 14		8	1908
6	174	9	2708
7	168	10	3008

11	2596	23	22
12	2086	24	6
13	1392	25	1
14	864	e = 22	
15	456	10	22202
16	218	11	42012
17	88	12	58742
18	30	13	63256
19	8	14	62396
20	2	15	54908
	e = 20	16	46352
9	6473	17	36112
10	10724	18	26906
11	13456	19	18756
12	13034	20	12456
13	11789	21	7468
14	9354	22	4472
15	7036	23	2408
16	4748	24	1208
17	3010	25	560
18	1728	26	238
19	914	27	88
20	426	28	30
21	187	29	8
22	68	30	2

PERCOLATION PERIMETER GROUPINGS

B — Square lattice

t = 4	g _{st} ($\times 1$)	6	54
1	1	7	22
t = 6		8	4
2	2		t = 11
t = 7		5	12
3	4	6	80
t = 8		7	136
3	2	8	80
4	9	9	28
5	1	10	4
t = 9			t = 12
4	8	5	2
5	20	6	60
6	4	7	252
t = 10		8	388
4	2	9	291
5	28	10	154

11	52	11	11772
12	9	12	12502
13	1	13	10480
t = 13		14	7508
6	16	15	4608
7	228	16	2406
8	777	17	1104
9	1152	18	396
10	986	19	124
11	644	20	28
12	325	21	4
13	112		
14	28		t = 16
15	4	7	2
t = 14		8	
6	2	9	2089
7	100	10	9750
8	818	11	24472
9	2444	12	38694
10	3676	13	44574
11	3530	14	41408
12	2644	15	33046
13	1660	16	23311
14	828	17	14385
15	332	18	8146
16	106	19	3982
17	22	20	1730
18	4	21	651
t = 15		22	
7	20	23	206
8	480	24	52
9	2804	25	9
10	7612		1

B — Honeycomb lattice

t = 3	g_{st} ($\times 1$)	t = 7	
1	1	5	15
t = 4		6	15
2	1.5	7	3
t = 5			t = 8
3	3	6	31.5
t = 6		7	60
4	7	8	37.5
5	3	9	12
6	0.5	10	1.5

t = 9		t = 12	
7	62	13	7078
8	177	14	11181
9	190	15	12937.5
10	111	16	11758
11	39	17	8895
12	9	18	5796
13	1	19	3258
t = 10		t = 19	
8	123	20	1522
9	471	21	565.5
10	744	22	164
11	705	23	37
12	449.5	24	6
13	207	t = 13	
14	69	11	0.5
15	15	12	1029
16	1.5	13	6927
t = 11		20160	
9	246	14	37635
10	1167	15	52311
11	2361	16	57960
12	3006	17	53949
13	2721	18	43728
14	1902	19	31536
15	1083	20	20355
16	492	21	11689
17	162	22	5889
18	33	23	2541
19	3	24	894
t = 12		234	
10	503	25	
11	2874	26	39
		27	3

B — Kagomé lattice

t = 4		g _{st} ($\times 1$)	t = 8	
1	t = 5	1	5	31
2	t = 6	2	6	12
3	t = 7	14/3 1/3	8	9
6			11	1
t = 7		t = 9		
4	12	6	$81 \frac{1}{3}$	
5	2	7	54	
7	2	9	$36 \frac{2}{3}$	
		10	11	

12	8	15	5952
15	2/3	16	3974
		17	5770
t = 10			
7	216	18	2190
8	220	19	1780
9	35	20	2318
10	133	21	468
11	88	22	748
13	45	23	674
14	10	25	258
16	9	26	126
19	1	28	70
		29	10
t = 11			
8	576	31	14
9	793	34	2
10	280		t = 14
11	454	11	11522
12	496	12	28524
13	58	13	30774
14	212	14	28078
15	138	15	38372
17	66	16	33712
18	19	17	21063
20	14	18	27850
23	2	19	18925
		20	10303
t = 12			
9	1550 $\frac{1}{3}$	21	14744
10	2724	22	6903
11	1576	23	3946
12	1620	24	6058
13	2344	25	1542
14	804	26	1729
15	877	27	1904
16	1020	28	139
17	153	29	639
18	371	30	430
19	270	32	195
21	118 $\frac{2}{3}$	33	68
22	46	35	49
24	29	38	9
27	4 $\frac{2}{3}$	41	1
30	1/3		t = 15
		12	317770 $\frac{2}{3}$
t = 13			
10	4210	13	89636
11	8940	14	117736
12	7432	15	120652 $\frac{2}{3}$
13	6536	16	151052
14	9726	17	162910

18	119496	19	637131
19	128790	20	608309
20	120932	21	645568
21	69820 $\frac{2}{3}$	22	472456
22	78394	23	418304
23	62468	24	425165
24	29484 $\frac{2}{3}$	25	247522
25	40976	26	237541
26	23430	27	216930
27	10837 $\frac{1}{3}$	28	99805
28	17521	29	122492
29	6096	30	86467
30	4432	31	36226
31	6050	32	55156
32	1082	33	27197
33	1763 $\frac{1}{3}$	34	12362
34	1689	35	21294
36	612	36	681
37	342	37	5339
39	179 $\frac{1}{3}$	38	6868
40	41	39	800
42	42 $\frac{2}{3}$	40	2052
45	8	41	1778
48	2/3	43	693
$t = 16$			
13	88129	44	360
14	279000	46	203
15	427488	47	32
16	500865	49	49
17	608253	52	9
18	716926	55	1

B — Triangular lattice (without holes)

$t = 6$		$t = 12$	
1	1	4	29
$t = 8$		5	21
2	3	6	14
$t = 9$		7	1
3	2	$t = 13$	
$t = 10$		5	66
3	9	6	42
4	3	7	30
$t = 11$		8	6
$t = 14$			
4	12	5	93
5	6	6	153

7	105		t = 19
8	69	8	5310
9	27	9	13488
10	3	10	18006
	t = 15	11	18732
6	298	12	14892
7	360	13	11310
8	264	14	7764
9	160	15	4614
10	86	16	2712
11	24	17	1374
12	2	18	606
	t = 16	19	198
6	306	20	42
7	840	21	6
8	918		t = 20
9	717	8	3408
10	444	9	20469
11	249	10	41673
12	117	11	51822
13	27	12	52992
14	3	13	45129
	t = 17	14	33981
7	1290	15	24900
8	2316	16	16188
9	2382	17	10023
10	1884	18	5676
11	1308	19	2879
12	720	20	1401
13	414	21	507
14	168	22	147
15	42	23	27
16	6	24	3
	t = 18		t = 21
7	1014	9	21372
8	4299	10	71644
9	6486	11	125448
10	6641	12	153614
11	5160	13	152658
12	3913	14	136014
13	2354	15	106416
14	1356	16	79446
15	729	17	55440
16	290	18	36576
17	87	19	22708
18	14	20	12912
19	1	21	7116

22	3442	25	128
23	1458	26	24
24	496	27	2

SIZE PERCOLATION GROUPINGS

C — Archimedean (3, 3, 4, 3, 4) (site problem)

s = 4	g _{st} ($\times 4$)	14	2926
5	4	15	6680
s = 2		16	8164
6	2	17	4632
7	8	18	848
s = 3		19	32
8	20	s = 9	
9	12	12	104
s = 4		13	612
8	2	14	2960
9	28	15	8780
10	66	16	20116
11	16	17	29908
s = 5		18	25312
9	8	19	9888
10	48	20	1266
11	180	21	36
12	156	s = 10	
13	20	12	24
s = 6		13	372
10	28	14	1998
11	108	15	9156
12	432	16	27284
13	676	17	62016
14	304	18	103726
15	24	19	110440
s = 7		20	68554
10	4	21	19204
11	72	22	1836
12	316	23	40
13	1092	s = 11	
14	2180	12	8
15	1928	13	128
16	528	14	1308
17	28	15	7104
s = 8		16	28436
11	20	17	86612
12	204	18	198268
13	988	19	350032

20	437312	17	91972
21	356728	18	279622
22	166292	19	645560
23	34668	20	1183662
24	2556	21	1639860
25	44	22	1606806
	s = 12	23	1032628
12	2	24	368306
13	40	25	59220
14	660	26	3452
15	4816	27	48
16	24572		

C — Archimedean (3, 3, 3, 4, 4) (site problem)

	s = 1	g _{st} ($\times 2$)	13	504
5		2	14	832
	s = 2		15	1084
6		2	16	464
7		2	17	110
8		1	s = 8	
	s = 3		11	4
7		2	12	120
8		4	13	504
9		10	14	1523
	s = 4		15	2750
8		2	16	3791
9		10	17	2906
10		33	18	1294
11		10	19	118
12		2	20	8
	s = 5		s = 9	
9		2	12	36
10		34	13	346
11		72	14	1512
12		68	15	4496
13		36	16	9004
	s = 6		17	13046
10		7	18	13130
11		90	19	8992
12		172	20	2300
13		254	21	314
14		254	s = 10	
15		36	12	5
16		4	13	158
	s = 7		14	1080
11		34	15	4758
12		204		

16	13604	23	55766
17	28980	24	9696
18	45462	25	854
19	52606	$s = 12$	
20	45831	13	4
21	21304	14	247
22	5518	15	2310
23	358	16	12693
24	16	17	48018
$s = 11$		18	137623
13	42	19	306124
14	608	20	541719
15	3670	21	755720
16	14956	22	828850
17	42802	23	676978
18	93434	24	392305
19	157478	25	123404
20	202272	26	21028
21	200886	27	1024
22	133464	28	32

CYCLOMATIC NUMBER DISTRIBUTIONS IN PERCOLATION

D — Triangular lattice

	$\Sigma_b bg_{sbt} (\times 1)$		
8	3	12	12
$s = 3$		13	330
9	6	14	1098
10	18	15	3198
$s = 4$		16	6504
10	15	17	8802
11	48	18	6084
12	87	$s = 8$	
$s = 5$		13	84
11	42	14	897
12	126	15	3420
13	324	16	10230
14	372	17	22494
$s = 6$		18	37251
12	126	19	41430
13	342	20	23856
14	1047	$s = 9$	
15	1746	14	432
16	1530	15	2478
		16	10962

17	32550	23	12435252
18	77244	24	19008123
19	142686	25	23432184
20	196716	26	22557609
21	187836	27	15089856
22	92496	28	5218950
$s = 10$		$s = 13$	
14	57	16	702
15	1548	17	10626
16	8256	18	67356
17	33522	19	322320
18	107844	20	1241400
19	258930	21	3838218
20	528585	22	10430136
21	822762	23	23661498
22	991227	24	46122252
23	826092	25	76516392
24	356391	26	105213888
$s = 11$		27	118095138
15	504	28	103613454
16	5088	29	63161082
17	28686	30	19879872
18	106518	$s = 14$	
19	354294	16	87
20	882774	17	4704
21	1886160	18	40776
22	3269094	19	254682
23	4492410	20	1112547
24	4797486	21	4240902
25	3559656	22	13165455
26	1366230	23	35992752
$s = 12$		24	83918157
15	48	25	169687884
16	2691	26	297846975
17	17442	27	446213892
18	97686	28	560840964
19	358230	29	578723592
20	1154871	30	467147454
21	3044358	31	261616854
22	6663264	32	75562266

D — Square matching site problem

$s = 2$	$\Sigma_b \text{bg}_{\text{sbt}} (\times 1)$	$s = 7$	
10	2	16	292
12	2	17	224
$s = 3$		18	1960
12	16	19	4120
14	16	20	7968
15	8	21	12092
16	4	22	21732
$s = 4$		23	22660
12	6	24	26520
14	78	25	24704
15	32	26	21956
16	96	27	13352
17	64	28	7820
18	72	29	3096
19	24	30	840
20	6	31	120
$s = 5$		32	12
14	64	16	100
15	24	17	72
16	332	18	1986
17	336	19	3872
18	568	20	10910
19	452	21	24892
20	612	22	51134
21	336	23	73640
22	208	24	119744
23	48	25	153072
24	8	26	177006
$s = 6$		27	172104
14	22	28	168700
16	396	29	127632
17	476	30	86776
18	1422	31	46364
19	2164	32	21414
20	3682	33	6580
21	3064	34	1400
22	4178	35	168
23	3264	36	14
24	2338		$s = 9$
25	1180	16	20
26	460	18	1360
27	80	19	2528
28	10	20	11440

21	32760	37	1800604
22	75996	38	879220
23	146044	39	341348
24	304820	40	101672
25	473628	41	21060
26	700344	42	3204
27	942580	43	288
28	1158604	44	18
29	1203972	s = 11	
30	1200980	18	208
31	1031344	19	372
32	810160	20	7300
33	519744	21	21692
34	294132	22	81332
35	135616	23	247844
36	49920	24	646804
37	12224	25	1492640
38	2176	26	3236880
39	224	27	6180136
40	16	28	11020012
s = 10		29	17637232
18	670	30	26672100
19	968	31	36246008
20	10406	32	45860536
21	31096	33	53549768
22	82660	34	57479148
23	219324	35	56049244
24	510996	36	50716588
25	948612	37	41178380
26	1795748	38	30463740
27	2903344	39	19695760
28	4335128	40	11292320
29	5806280	41	5528144
30	7281978	42	2318500
31	8178288	43	756984
32	8407562	44	190168
33	7738904	45	33960
34	6627886	46	4520
35	4895912	47	360
36	3241420	48	20

D — Kagomé lattice site problem

	$s = 2$	$\Sigma_b \text{bg}_{\text{stb}} (\times 3)$	12	71826
5		6	13	296724
	$s = 3$		14	1141584
6		30	15	1241916
	$s = 4$			$s = 13$
7		120	10	2376
	$s = 5$		11	2682
7		24	12	113640
8		426	13	306876
	$s = 6$		14	1349736
6		6	15	3912288
8		204	16	3762834
9		1416		$s = 14$
	$s = 7$		10	546
7		48	11	12126
9		1140	12	40842
10		4548	13	511212
	$s = 8$		14	1397292
8		258	15	5612268
10		5496	16	13185468
11		14220	17	11375370
	$s = 9$			$s = 15$
9		1206	9	42
10		936	11	8190
11		23052	12	53886
12		43860	13	329904
	$s = 10$		14	2159004
9		378	15	6383382
10		4932	16	21974364
11		8880	17	43892388
12		88956	18	34321086
13		134250		$s = 16$
	$s = 11$		10	612
8		42	12	65484
10		3504	13	255342
11		18636	14	2012856
12		56616	15	9002940
13		325044	16	28199694
14		408894	17	82526616
	$s = 12$		18	144709680
9		384	19	103371816
11		21972		$s = 17$
			11	4824
			12	10158

13	396048	17	118643394
14	1397940	18	301417152
15	10394394	19	473127060
16	38126082	20	310901274

D — Honeycomb lattice site problem

s = 2	$\Sigma_b \text{bg}_{\text{stb}} (\times 2)$	s = 12	
4	3	9	234
s = 3		10	10797
5	12	11	68670
s = 4		12	156972
6	42	13	152394
s = 5		14	46530
6	24	s = 13	
7	120	9	30
s = 6		10	5616
6	6	11	69870
7	150	12	275772
8	315	13	486714
s = 7		14	395352
7	42	15	104496
8	720	s = 14	
9	744	10	2097
s = 8		11	54420
8	552	12	354312
9	2478	13	998700
10	1722	14	1457295
s = 9		15	1001676
8	216	16	234312
9	3126	s = 15	
10	7536	10	510
11	3936	11	34248
s = 10		12	355644
8	33	13	1524900
9	2166	14	3428742
10	13623	15	4197072
11	21006	16	2500512
12	9054	17	524244
s = 11		s = 16	
9	882	10	57
10	14862	11	17130
11	47766	12	295326
12	57480	13	1849878
13	20580	14	6028185
		15	11161506

16	11757504		s = 19
17	6162480	11	138
18	1171950	12	64644
	s = 17	13	1285596
11	6198	14	10333392
12	209940	15	47614188
13	1876770	16	137698308
14	8488092	17	258202302
15	22152534	18	311675004
16	34968960	19	227833902
17	32137488	20	87704424
18	15050304	21	12997152
19	2616288		s = 20
	s = 18	12	25860
11	1386	13	892020
12	128046	14	9492018
13	1650396	15	55892538
14	10041624	16	206922024
15	35364792	17	506625396
16	77356395	18	832200162
17	105800778	19	896996238
18	86255844	20	593959458
19	36453678	21	209710524
20	5835012	22	28922142

D — Archimedean lattice (3, 3, 4, 3, 4) site problem

s = 2	Σ_b bg _{stb} ($\times 4$)	12	2652
6	2	13	3680
7	8	14	1560
	s = 3	15	120
8	44		s = 7
9	24	10	36
	s = 4	11	656
8	8	12	2644
9	104	13	8220
10	210	14	15100
11	48	15	12272
	s = 5	16	3240
9	48	17	168
10	248		s = 8
11	832	11	224
12	644	12	2190
13	80	13	9696
	s = 6	14	26442
10	208	15	55872
11	724	16	63322

17	33960	14	21148
18	6056	15	107820
19	224	16	407184
s = 9		17	1168464
12	1344	18	2518668
13	7436	19	4211012
14	33664	20	5012660
15	92464	21	3902412
16	196812	22	1753036
17	275136	23	357472
18	218848	24	25948
19	82256	25	440
20	10396	s = 12	
21	288	12	38
s = 10		13	768
12	356	14	12156
13	5408	15	84764
14	27408	16	409734
15	117676	17	1453032
16	327452	18	4178896
17	697796	19	9142376
18	1100140	20	15939366
19	1110356	21	21095484
20	656814	22	19821760
21	178864	23	12241136
22	16796	24	4242200
23	360	25	669848
s = 11		26	38508
12	136	27	528

D — Archimedean lattice (3, 3, 3, 4, 4) site problem

s = 2	$\Sigma_b \text{bg}_{\text{stb}} (\times 2)$	s = 5	
6	2	9	14
7	2	10	184
8	1	11	318
s = 3		12	280
7	6	13	144
8	8	s = 6	
9	20	10	56
s = 4		11	626
8	10	12	1028
9	38	13	1378
10	102	14	1292
11	30	15	180
12	6	16	20

s = 7		19	524380
11	322	20	435309
12	1724	21	196808
13	3826	22	50046
14	5708	23	3222
15	6828	24	144
16	2832	s = 11	
17	660	13	754
s = 8		14	10084
11	48	15	56894
12	1322	16	216292
13	5008	17	579562
14	13914	18	1189958
15	22892	19	1893344
16	29028	20	2301600
17	21208	21	2177084
18	9158	22	1396480
19	826	23	569130
20	56	24	97684
s = 9		25	8540
12	488	s = 12	
13	4316	13	84
14	17404	14	4770
15	47576	15	41758
16	88356	16	215426
17	118700	17	764342
18	112864	18	2066141
19	74098	19	4346780
20	18600	20	7293506
21	2512	21	9678184
s = 10		22	10131540
12	81	23	7961512
13	2390	24	4476120
14	15100	25	1382338
15	61894	26	232628
16	163962	27	11264
17	326878	28	352
18	480335		

D — Square bond problem

s = 1	$\sum_s s g_{stb} (\times 2)$	10	72
6	4		
s = 2		s = 4	
8	18	8	4
s = 3		11	160
9	16	12	275

	s = 5		19	97920
10	40		20	58257
12	180			
13	960		13	480
14	1044		14	1072
	s = 6		15	1476
11	84		16	24984
12	240		17	42516
14	2324		18	84992
15	4704		19	273920
16	3990		20	460040
	s = 7		21	432200
10	12		22	222020
13	1092			
14	1176		12	48
15	2688		14	558
16	16240		15	6552
17	21696		16	5904
18	15264		17	53080
	s = 8		18	188160
12	154		19	281096
14	1824		20	809612
15	7664		21	1636008
16	7144		22	2281884
17	36576		23	1867624
18	89748		24	845746

D — Honeycomb bond problem

	s = 1	$\Sigma_s sg_{stb} (\times 1)$	10	5184
4		6		
	s = 2		8	216
5		18	10	5940
	s = 3		11	15552
6		56		
	s = 4		9	990
7		180	10	1050
	s = 5		11	23940
7		36	12	46510
8		558		
	s = 6		9	330
6		6	10	3990
8		252	11	9240
9		1708	12	89892
	s = 7		13	138930
7		42		
9		1296	8	30
	s = 11			

10	2904		s = 15
11	14982	9	26
12	56736	11	5796
13	321840	12	36834
14	414792	13	267840
	s = 12	14	1730016
9	264	15	5652390
11	17856	16	20519424
12	58712	17	41415840
13	289848	18	33004544
	s = 13		s = 16
14	1112436	10	378
15	1239056	12	45900
	s = 14	13	180702
10	1620	14	1618176
11	2262	15	7296948
12	91416	16	25115160
13	260268	17	76400448
14	1292508	18	135462936
15	3764712	19	98498952
16	3701418		s = 17
	s = 15	11	2970
10	390	12	7344
11	8268	13	276960
12	33768	14	1037808
13	408492	15	8308410
14	1218804	16	31418934
15	5298120	17	105743982
16	12555000	18	277089606
17	11054610	19	439782480
		20	293866272

D — Kagomé lattice site problem

s = 2	$\sum_b b^2 g_{sbt} (\times 3)$		s = 7
5	6	7	384
s = 3		9	8088
6	66	10	32232
s = 4			s = 8
7	408	8	2472
s = 5		10	46104
7	96	11	117972
8	1980		s = 9
s = 6		9	13266
6	36	10	8406
8	1164	11	223344
9	8316	12	416466

	s = 10		15	89507916
9	4344		16	208467996
10	61158		17	176236950
11	94416			s = 15
12	973692		9	882
13	1435470		11	162306
	s = 11		12	1105230
8	588		13	6111624
10	46632		14	40617600
11	255828		15	113174922
12	681552		16	377802300
13	3958824		17	746740380
14	4861974		18	572993958
	s = 12			s = 16
9	6144		10	13878
11	325416		12	1403460
12	1066932		13	5488974
13	3967332		14	40162224
14	15296832		15	179523264
15	16256316		16	531568206
	s = 13		17	1520199504
10	41856		18	2636009496
11	41490		19	1850067696
12	1841832			s = 17
13	4835928		11	117576
14	19813620		12	225234
15	57146988		13	9075648
16	53773968		14	31074372
	s = 14		15	221592582
10	9954		16	799876968
11	231498		17	2375467950
12	693798		18	5918963946
13	8980896		19	9187911360
14	23326068		20	5937914886

FIXED SIZE ENERGY GROUPINGS

E — Honeycomb lattice

	s = 1	g _{se} ($\times 2$)		s = 5	
3		2		7	36
	s = 2				s = 6
4		3		6	1
	s = 3			8	93
5		6			s = 7
	s = 4			7	6
6		14		9	244

	s = 8		s = 16	
8	27		10	3
10	648		12	1113
	s = 9		14	36405
9	110		16	428955
11	1728		18	2062784
	s = 10			s = 17
8	3		11	42
10	399		13	5340
12	4651		15	127170
	s = 11		17	1318254
9	24		19	5794056
11	1362			s = 18
13	12630		12	356
	s = 12		14	22941
10	135		16	433941
12	4468		18	4035356
14	34566		20	16325904
	s = 13			s = 19
9	2		11	6
11	636		13	2244
13	14244		15	91622
15	95312		17	1456674
	s = 14		19	12310578
10	27		21	46133390
12	2631			s = 20
14	44706		12	87
16	264387		14	11838
	s = 15		16	348120
11	198		18	4830054
13	10050		20	37445847
15	138938		22	130703016

E — Square lattice

	s = 1	$g_{se} (\times 1)$	12	55
4		1		
	s = 2		10	2
6		2	12	40
	s = 3		14	174
8		6		
	s = 4		12	22
8		1	14	168
10		18	16	570
	s = 5			
10		8	12	6

14	134	18	864
16	677	20	9354
18	1908	22	62396
s = 9		24	297262
12	1	26	1056608
14	72	28	2448760
16	656	30	3329608
18	2708		s = 15
20	6473	16	6
s = 10		18	456
14	30	20	7036
16	482	22	54908
18	3008	24	317722
20	10724	26	1359512
22	22202	28	4401192
s = 11		30	9436252
14	8	32	11817582
16	310		s = 16
18	2596	16	1
20	13456	18	218
22	42012	20	4748
24	76886	22	46352
s = 12		24	303068
14	2	26	1563218
16	151	28	6095764
18	2086	30	18173796
20	13034	32	36285432
22	58742	34	42120340
24	163494		s = 17
26	268352	18	88
s = 13		20	3010
16	68	22	36112
18	1392	24	276464
20	11789	26	1603984
22	63256	28	7477928
24	250986	30	26922156
26	633748	32	74496544
28	942651	34	139297108
s = 14		36	150682450
16	22		

E — Triangular lattice

s = 1	$g_{se} (\times 1)$	s = 3	
6	1	12	2
s = 2			
10	3	14	9

s = 4		30	4817
14	3	32	12453
16	6	34	28305
18	29	36	55411
s = 5		38	92205
16	6	40	101679
18	21	42	65822
20	60	s = 11	
22	99	24	24
s = 6		26	249
18	14	28	1164
20	42	30	3876
22	129	32	12264
24	281	34	34263
26	348	36	78990
s = 7		38	165144
18	1	40	295356
20	30	42	438264
22	105	44	434784
24	276	46	251655
26	732	s = 12	
28	1248	24	2
30	1260	26	117
s = 8		28	702
20	6	30	3163
22	69	32	10557
24	246	34	32901
26	720	36	91389
28	1737	38	226125
30	3795	40	482850
32	5472	42	929414
34	4644	44	1531383
s = 9		46	2050899
22	27	48	1852892
24	160	50	969819
26	609	s = 13	
28	1818	26	27
30	4697	28	414
32	9990	30	2126
34	19014	32	8418
36	23662	34	29601
38	17382	36	89680
s = 10		38	250641
22	3	40	631218
24	86	42	1431751
26	432	44	2845248
28	1458	46	5093199

48	7761168	38	252177
50	9484524	40	694122
52	7876554	42	1782600
54	3762517	44	4157097
s = 14		46	8736174
26	3	48	16309377
28	168	50	27275403
30	1320	52	38620725
32	6336	54	43453965
34	23721	56	33417534
36	81183	58	14680890

E — Square covering site problem

s = 1	$g_{se} (\times 2)$	24	344
6	2	26	1924
s = 2		28	4035
10	6	30	10858
s = 3		32	13259
12	4	34	2958
14	18	s = 9	
s = 4		22	26
12	1	24	168
16	37	26	1076
18	50	28	3336
s = 5		30	13512
16	12	32	25240
18	26	34	56634
20	192	36	47320
22	142	38	8134
s = 6		s = 10	
18	16	22	4
20	102	24	58
22	246	26	580
24	874	28	2266
26	390	30	9360
s = 7		32	29444
18	2	34	83758
20	24	36	152964
22	226	38	266710
24	640	40	163340
26	1826	42	22050
28	3508	s = 11	
30	1086	24	16
s = 8		26	236
20	6	28	1372
22	82	30	6260

32	23720	46	4832436
34	74098	48	1784168
36	222372	50	162466
38	493862	$s = 13$	
40	869372	26	12
42	1169136	28	316
44	545580	30	2122
46	60146	32	11448
$s = 12$		34	42428
24	1	36	170104
26	76	38	520236
28	743	40	1612728
30	3704	42	3966084
32	17174	44	9399652
34	58860	46	15765404
36	208354	48	23066864
38	553224	50	19101104
40	1507761	52	5711504
42	2822608	54	440750
44	4625299		

REFERENCES

- [1] SYKES M. F., WATTS M. G., GAUNT D. S., *J. Phys.*, **A8**, 1448 (1975).
- [2] SYKES M. F., GAUNT D. S., MARTIN J., MATTINGLY S. R., ESSAM J. W. *J. Math. Phys.*, **14**, 1071 (1973).
- [3] DUARTE J. A. M. S., MARQUES M. C. A. M., *Nuovo Cim.* **B54**, 508 (1979).
- [4] SYKES M. F., GAUNT D. S., GLEN M. *J. Phys.*, **A9**, 715 (1976).
- [5] HILEY B. J., SYKES M. F., *J. Chem. Phys.*, **34**, 1531 (1961).
- [6] ENTING I. G., *J. Phys.*, **A13**, 3713 (1980).
- [7] BLEASE J., ESSAM J. W., PLACE C. M., *J. Phys.*, **C11**, 4009 (1978).
- [8] SYKES M. F., GLEN M., *J. Phys.*, **A9**, 87 (1976).
- GAUNT D. S., SYKES M. F., RUSKIN H. J., *J. Phys.*, **A9**, 1899 (1976).
- GAUNT D. S., RUSKIN H. J., *J. Phys.*, **A11**, 1369 (1978).
- [9] PETERS H. P., STAUFFER D., HOLTERS H. P., LOEWENICH L., *Zeit f. Phys.* **B35**, 399 (1979).
- [10] SYKES M. F., GAUNT D. S., GLEN M., *J. Phys.*, **A14**, 287 (1981).
- [11] DUARTE J. A. M. S., *J. Physique*, **40**, 845 (1979).
- [12] CHERRY R. J., PH. D. THESIS, London, unpublished (1979).
- [13] GAUNT D. S., MIDDLEMISS K. M., TORRIE G., WHITTINGTON S. G., *J. Phys.*, **A13**, 3029 (1980).
- [14] DUNN A. G., ESSAM J. W., RITCHIE D. S., *J. Phys.*, **C8**, 4219 (1975).
- [15] AGRAWAL P., REDNER S., REYNOLDS P. J., STANLEY H. E., *J. Phys.*, **A12**, 2073 (1979).
- [16] DUARTE J. A. M. S., RUSKIN, H. J., *Zeit. f. Phys.*, **B**, in press (1981).
- [17] LUBENSKY T. C., ISAACSON, J., *Phys. Rev.*, **A20**, 2130 (1979).
- [18] HARRIS A. B., LUBENSKY T. C., *Phys. Rev.*, **B23**, 3591 (1981).
- [19] DUARTE J. A. M. S., RUSKIN H. J., *J. Physique* in press (1981).